

Victoria University of Wellington

Te Whare Wānanga o te Ūpoko o te Ika a Maui



Sparsity of the Hawking Flux

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1 Introduction and History

2 Sparsity Results

- Semi-Analytic Estimates
- Numerical Estimates
- Caveat Super-Radiance

3 Conclusion and Outlook

Black Holes Radiate Differently

- Black hole radiation usually compared to black body radiation
- However: Fundamental differences!
- Sparsity: 40-year-old result¹
- Different view points:
 - $\lambda_{\text{thermal}} \gg r_H$
 - Long lifetime of black hole
 - $\tau_{\text{loc}} \ll \tau_{\text{gap}} \leftarrow \text{our focus}$

¹E.g.: D. Page, PRD **13**(1976) 198; PRD **14**(1976) 3260; PRD **16**(1977) 2404

Measures of Sparsity

- Natural time-scales set by black body spectrum & emission probability

$$d\Gamma = \frac{g}{(2\pi)^3} \frac{c (\hat{k} \cdot \hat{n})}{\exp\left((\epsilon(\vec{k}) - \mu)/k_B T\right) + s} d^3\vec{k} dA, \quad \tau_{\text{gap}} := \frac{1}{\Gamma}$$

- Many obvious choices, e.g.:

-

$$\eta_{\text{peak number}} = \frac{\tau_{\text{gap}}}{\tau_{\text{peak number}}}$$

-

$$\eta_{\text{peak energy}} = \frac{\tau_{\text{gap}}}{\tau_{\text{peak energy}}}$$

-

$$\eta_{\text{average energy}} = \frac{\tau_{\text{gap}}}{\tau_{\text{average energy}}}$$

- Bolometric measure: $\eta_{\text{binned}} = \frac{1}{\int_0^\infty \frac{2\pi}{ck} \frac{d\Gamma}{dk} dk}$

What we look at

- Different types of black holes & particles in ray optics limit
- Black body approximation: Semi-analytic results
- Grey bodies: Numerical results
- Neglecting:
 - **Adiabaticity constraints:** Backreaction and how fast the space-time may change
→ *infrared cutoff*
 - **Phase space constraints:** Photon energy cannot exceed BH mass
→ *ultraviolet cutoff*
- Aiming for conservative estimates

Schwarzschild BH, Semi-analytic

For massless particles:

Type	$\eta_{\text{peak number}}$	$\eta_{\text{peak energy}}$	$\eta_{\text{avg. frequency}}$	η_{binned}
Bosons \approx	$\frac{32\pi^2(2+W(-2e^{-2}))}{27g\zeta(3)}$ 15.508/ g	$\frac{32\pi^2(3+W(-3e^{-3}))}{27g\zeta(3)}$ 27.465/ g	$\frac{16\pi^6}{405g\zeta(3)^2}$ 26.285/ g	$\frac{128}{9g}$ 14.222/ g
Fermions \approx	$\frac{128\pi^2(2+W(2e^{-2}))}{81g\zeta(3)}$ 28.773/ g	$\frac{128\pi^2(3+W(3e^{-3}))}{81g\zeta(3)}$ 40.624/ g	$\frac{224\pi^6}{3645g\zeta(3)^2}$ 40.88/ g	$\frac{256}{9g}$ 28.444/ g
Boltzmann \approx	$\frac{64\pi^2}{27g}$ 23.395/ g	$\frac{32\pi^2}{9g}$ 35.092/ g	$\frac{32\pi^2}{9g}$ 35.092/ g	$\frac{64\pi^2}{27g}$ 23.395/ g

Things can be summarised, e.g.:

$$\eta_{\text{peak number}} = \frac{\lambda_{\text{thermal}}^2}{A_{\text{H}}} \frac{2 + W(2s e^{-\mu-2})}{2\text{Li}_3(-s e^{-\mu})/(-s)}, \quad s \in \{-1, 0, 1\}$$

Generalising Semi-Analytics

- For Reissner-Nordström & Kerr in absence of super-radiance:

$$\eta = \eta_{\text{Schwarzschild}} \times \frac{r_+^2}{(r_+ - r_-)^2} \geq \eta_{\text{Schwarzschild}}$$

- Dirty black holes:

$$\eta = \eta_{\text{Schwarzschild}} \times \frac{e^{\phi_H}}{1 - 8\pi G_N \rho_H r_H^2 / c^4}, \quad \phi_H = \frac{4\pi G_N}{c^4} \int_{r_H}^{\infty} \frac{(\rho - p_r) r}{1 - 2m(r)/r} dr \geq 0$$

- Particle rest mass:

$$\frac{\int_0^\infty \frac{x^2}{\exp(\sqrt{z^2 + x^2}) \mp 1} dx}{\int_0^\infty \frac{x^2}{\exp(x) \mp 1} dx} \leq 1 \quad \implies \quad \eta \geq \eta_{\text{Schwarzschild}}$$

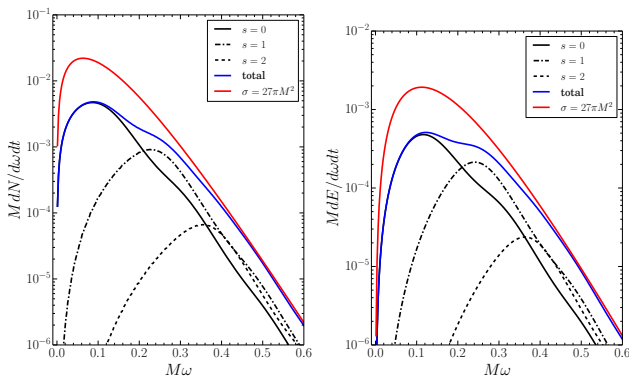
- Extending to graviton results for Tangherlini metric²:

$$\eta_{\text{Tang,grav}} = \frac{1}{(d-3)!2\pi} \frac{\Gamma(\frac{d-2}{2})}{\pi^{(d-4)/2}} \frac{(d-2 + W((d-2)se^{-\mu-d+2}))}{\frac{\text{Li}_{d-1}(-se^{-\mu})}{(-s)}} \frac{\lambda_{\text{thermal}}^{d-2}}{gA_H} \stackrel{d \lesssim 10}{>} 1$$

²Unpublished, compare to S. Hod, Phys. Lett. B, **756**, 133-136 (2016)

Numerical Results

Ray optics limit \rightarrow Turn Regge-Wheeler equation into Shabat-Zakharov system, use product calculus³



Grey body factors make things even sparser!

³F. Gray, M. Visser, arXiv:1512.05018, consistent with Teukolsky & Page

Super-Radiance

- q and $L \rightarrow$ chemical potentials/ergo-regions
- For $\omega < \omega_{\text{crit}}$: $G_{\ell,m}$ and $\langle N \rangle_\epsilon$ now $< 1 \Rightarrow$ Super-radiance
- Both effects act in parallel
- Therefore (here for angular momentum and binned η):

$$\frac{1}{\eta_{\text{super-radiance}}} + \frac{1}{\eta_{\text{Hawking}}} = 2\pi \sum_{\ell,m} \int_0^{\omega_{\text{crit}}} G_{\ell,m} N_\epsilon \frac{d\omega}{\omega} + 2\pi \sum_{\ell,m} \int_{\omega_{\text{crit}}}^{\infty} G_{\ell,m} N_\epsilon \frac{d\omega}{\omega}$$

- For $\kappa \rightarrow 0$, super-radiance dominates:

$$\eta_{\text{super-radiant}} = \mathcal{O}(\epsilon^0); \quad \eta_{\text{Hawking}} = \mathcal{O}(\epsilon^{2s+2}) \gggg 1$$

$$\text{since } \epsilon := \hbar\omega_{\text{crit}}/(k_B T_H) \gggg 1$$

Conclusion and Outlook

Summary

Even conservative semi-analytical estimates are lower bounds for η_{Hawking} !

- $1 \ll \eta_{\text{semi-analytic, Schwarzschild}} \leq \eta_{\text{semi-analytic, general}}$
- After some care with g : $\eta_{\text{Boson}} < \eta_{\text{Boltzmann}} < \eta_{\text{Fermion}}$
- \implies **Sparse flux!**
- Unlike black/grey bodies, Hawking radiation is a long 2-body decay chain

Outlook

- Analogue black holes?
- Correlations in interstitial gaps?

Thank you!

Questions?

- Wave equation on curved background \rightarrow stationary, 1D-Schrödinger equation in r^* with Regge-Wheeler potential

$$V(r^*) = \left[1 - \frac{2M}{r(r^*)} \right] \left[\frac{\ell(\ell+1)}{r(r^*)^2} + \frac{(1-s^2)2M}{r(r^*)^3} \right]$$

- Rephrase Schrödinger eqn. as Shabat-Zakharov system
- Use (numerical) product calculus to solve for Bogoliubov coefficients (F. Gray, M. Visser, arXiv:1512.05018):

$$\begin{pmatrix} \alpha & \beta^* \\ \beta & \alpha^* \end{pmatrix} = \prod_{-\infty}^{\infty} \left(\mathcal{I} - \frac{i}{4x} V(u^*) \begin{pmatrix} 1 & \exp(-4ixu^*) \\ -\exp(4ixu^*) & -1 \end{pmatrix} \right) du^*$$

with $u = r/2M$, $u^* = r^*/2M$

- Check consistency with earlier results by Teukolsky & Page: ✓

Shabat-Zakharov System

- Start with a wave equation on curved background:

$$\nabla_\mu \nabla^\mu \psi = 0$$

- Rephrase resulting Regge-Wheeler equation with the ansatz

$$\psi(x) = a(x) \frac{\exp(+i\varphi)}{\sqrt{\varphi'}} + b(x) \frac{\exp(-i\varphi)}{\sqrt{\varphi'}}$$

and gauge condition

$$\frac{d}{dx} \left(\frac{a}{\sqrt{\varphi'}} \right) e^{+i\varphi} + \frac{d}{dx} \left(\frac{b}{\sqrt{\varphi'}} \right) e^{-i\varphi} = 0$$

- With $\omega(x)^2 := E - V(x)$ and $\rho := \varphi'' + i(\omega(x)^2 - (\varphi')^2)$ get S-Z equations:

$$\frac{d}{dx} \begin{pmatrix} a(x) \\ b(x) \end{pmatrix} = \frac{1}{2\varphi'} \begin{pmatrix} i \operatorname{Im}(\rho) & \rho \exp(-2i\varphi) \\ \rho^* \exp(2i\varphi) & -i \operatorname{Im}(\rho) \end{pmatrix} \begin{pmatrix} a(x) \\ b(x) \end{pmatrix}$$

- Formally solve S-Z equations by path-ordered integration

$$\begin{pmatrix} \alpha & \beta^* \\ \beta & \alpha^* \end{pmatrix} = \mathcal{P} \exp \left(-\frac{i}{4x} \int_{-\infty}^{\infty} V(u^*) \begin{pmatrix} 1 & \exp(-4ixu^*) \\ -\exp(4ixu^*) & -1 \end{pmatrix} du^* \right)$$

- Equivalent to product calculus:

$$\int_a^b f(x) dx = \lim_{N \rightarrow \infty} \sum_{k=0}^{N-1} f(t_k) \Delta x_k$$

$$\Leftrightarrow \prod_a^b (\mathcal{I} + A(x) dx) \equiv \lim_{N \rightarrow \infty} \prod_{k=0}^{N-1} (\mathcal{I} + A(x_k^*) \Delta x_k)$$

- Many approaches to numerical evaluation, including:
 - Peano-Baker series
 - Simpson rule
 - Truncating the defining limits
 - ...