

Causal Nature and Dynamics of Trapping Horizons in Black Hole Collapse

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Reference:

[arXiv: 1601.05109v2](#)

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Introduction

- **Spherically symmetric** metric in comoving coordinates with t “cosmic time”:

$$ds^2 = a^2(r, t)dt^2 + b^2(r, t)dr^2 + R^2(r, t)d\Omega^2$$

- **Proper time** and **proper distance** operators:

$$D_t \equiv \frac{1}{a} \frac{\partial}{\partial t} \Rightarrow U \equiv D_t R \qquad D_r \equiv \frac{1}{b} \frac{\partial}{\partial r} \Rightarrow \Gamma \equiv D_r R$$

- **Perfect Fluid:** $T^{\mu\nu} = (e + p)u^\mu u^\nu + pg^{\mu\nu}$

- Constraint equation (integrating G_{00}) :

$$\Gamma^2 = 1 + U^2 - \frac{2M}{R}$$

- Mister-Sharp Mass : $M = \int 4\pi e R^2 dR$

$$D_t M = -4\pi p R^2 U$$

Trapping Horizons

Expansion of **ingoing/outgoing** null-rays :

$$k^a / l^a = \left(\frac{1}{a}, \pm \frac{1}{b}, 0, 0 \right) \implies \theta_{\pm} = h^{cd} \nabla_c k_d = \frac{2}{R} (U \pm \Gamma)$$
$$h_{ab} = g_{ab} + \frac{1}{2} (k_a l_b + l_a k_b) \qquad k^a l_a = -2$$

Black Hole / Cosmological horizon : $\theta_{\pm} = 0 \implies \frac{1}{a} \frac{dR}{dt} \Big|_{\pm} \implies \Gamma^2 = U^2$

$$R = 2M$$

The horizon condition is independent of the slicing and holds also within a non-vacuum moving medium

The so-called **apparent horizon** of a black hole (which is a future trapping horizon) is the **outermost trapped surface for outgoing radial null rays** while the **trapping horizon** for an expanding universe (which is a past trapping horizon) is **foliated by the innermost anti-trapped surfaces for ingoing radial null rays**.

Causal Nature

$$\alpha \equiv \frac{\mathcal{L}_v \theta_v}{\mathcal{L}_{nv} \theta_v} \left\{ \begin{array}{l} \alpha > 0 : \text{space-like} \\ \alpha = 0 / \infty : \text{null} \\ \alpha < 0 : \text{time-like} \end{array} \right.$$

$$\text{Lie Derivatives: } \left\{ \begin{array}{l} \mathcal{L}_+ \theta_v = \mathcal{L}_k \theta_v = k^a \partial_a \theta_v = \left(\frac{1}{a} \frac{\partial}{\partial t} + \frac{1}{b} \frac{\partial}{\partial r} \right) \theta_v \\ \mathcal{L}_- \theta_v = \mathcal{L}_l \theta_v = l^a \partial_a \theta_v = \left(\frac{1}{a} \frac{\partial}{\partial t} - \frac{1}{b} \frac{\partial}{\partial r} \right) \theta_v \end{array} \right.$$

$$\mathcal{L}_{\pm} \theta_v = (D_t \pm D_r) \theta_v$$

$$\alpha = \left. \frac{4\pi R^2 (e + p)}{1 - 4\pi R^2 (e - p)} \right|_H$$

Horizon Velocity

3-velocity of the horizon with respect the matter: $v_H \equiv \left(\frac{b}{a} \frac{dr}{dt} \right)_H$

$$\theta_v = 0 \quad \Rightarrow \quad D_t \theta_v + \frac{b}{a} \frac{dr}{dt} D_r \theta_v = 0$$

$$v_H \equiv - \frac{D_t \theta_v}{D_r \theta_v} \quad \Rightarrow \quad v_H = - \frac{D_t (\Gamma^2 - U^2)}{D_r (\Gamma^2 - U^2)} \Big|_H$$

$$v_H = - \frac{\mathcal{L}_+ \theta_v + \mathcal{L}_- \theta_v}{\mathcal{L}_+ \theta_v - \mathcal{L}_- \theta_v} \Big|_H \quad \Rightarrow \quad v_H = \pm \frac{1 + \alpha}{1 - \alpha}$$

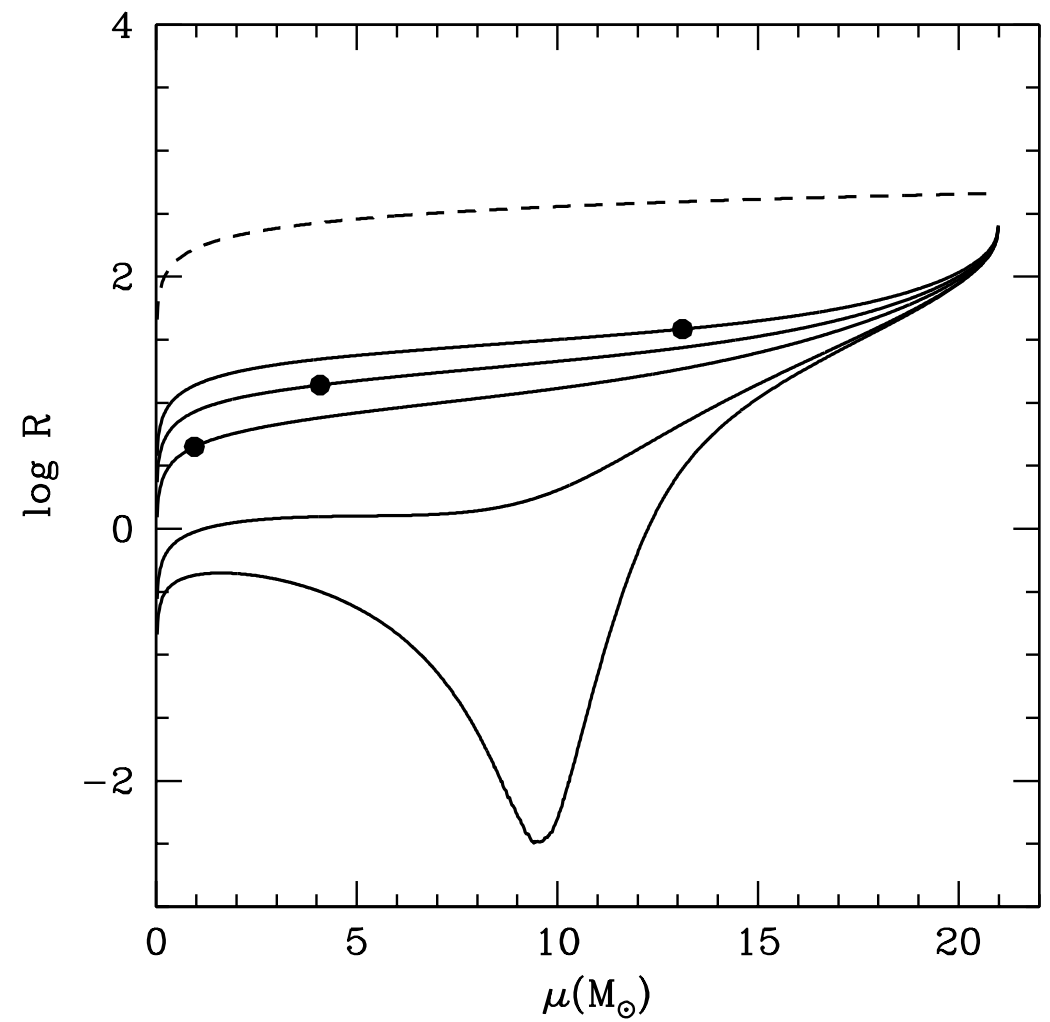
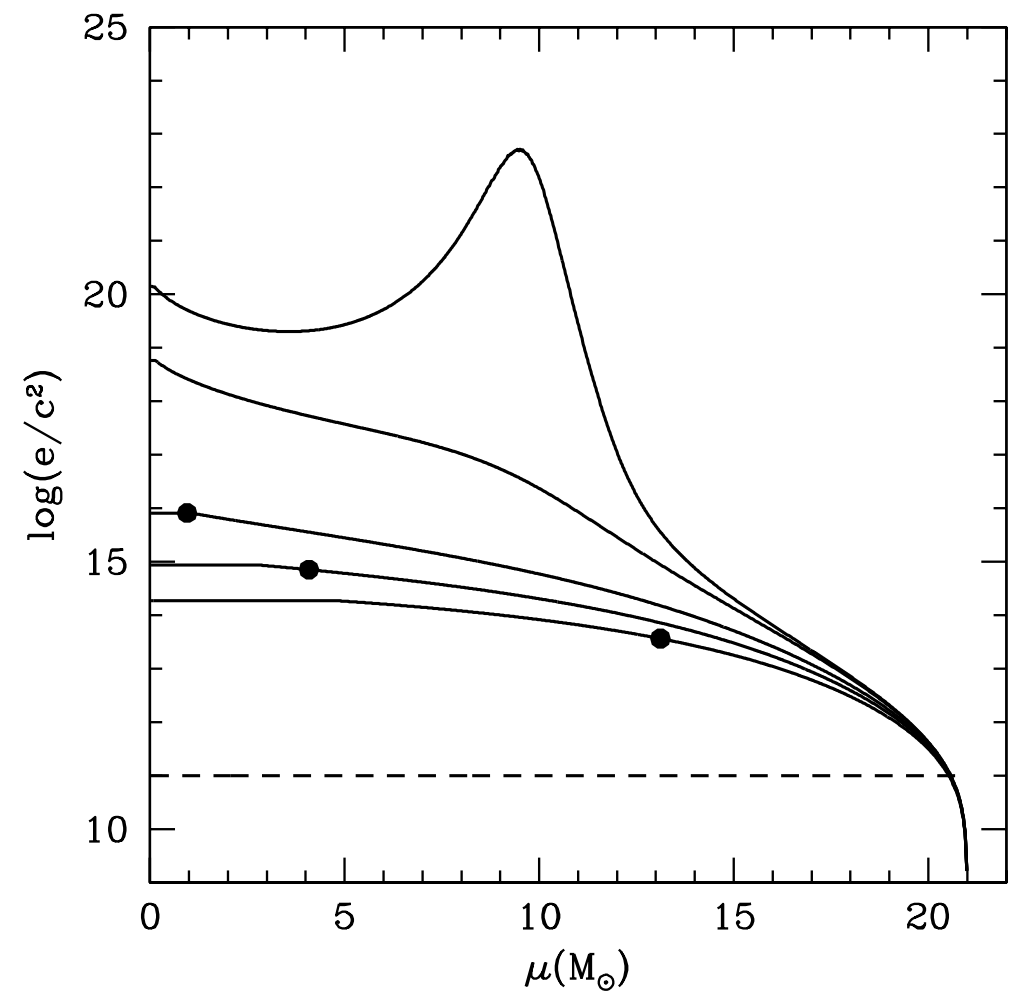
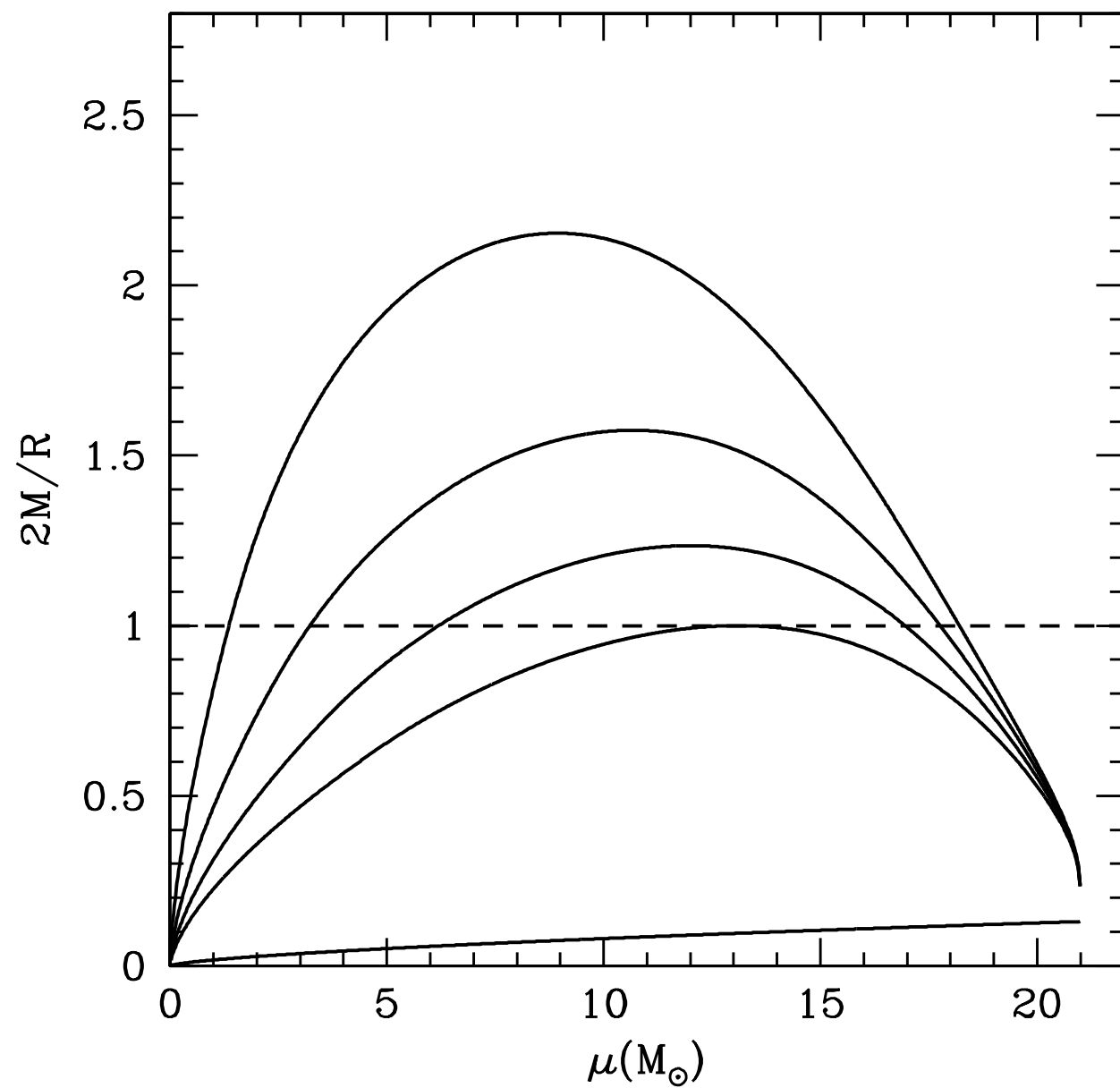
$$v_H = - \frac{U}{\Gamma} \Big|_H \frac{1 + 8\pi R^2 p}{1 - 8\pi R^2 e} \Big|_H$$

$$\left\{ \begin{array}{l} |v_H| > 1: \text{ space-like} \\ |v_H| = 1: \text{ null} \\ |v_H| < 1: \text{ time-like} \end{array} \right.$$

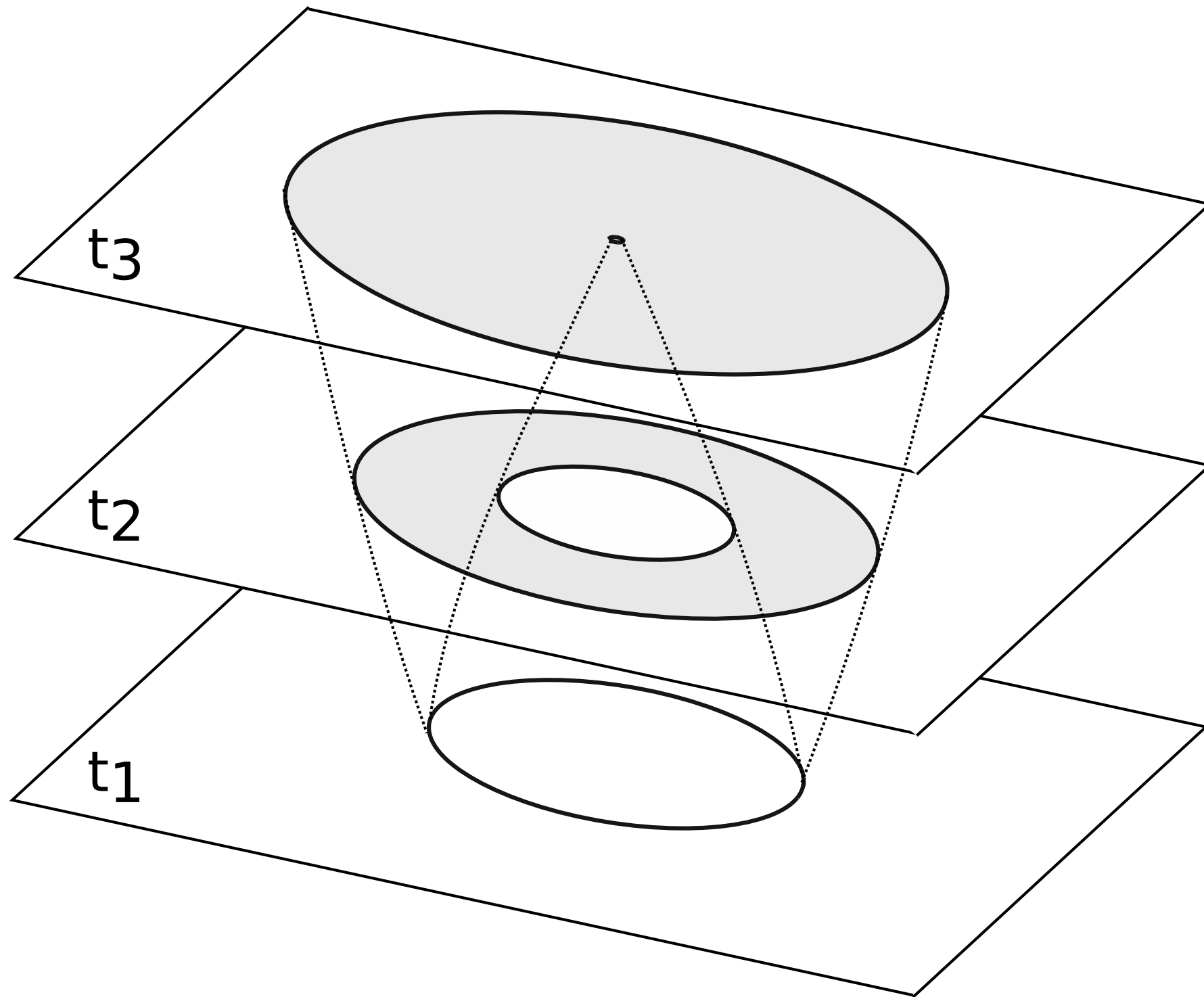
May & White (1966)

Homogenous initial density

$$p = K \rho^\gamma \quad (\gamma = 5/3, K = 0.05)$$



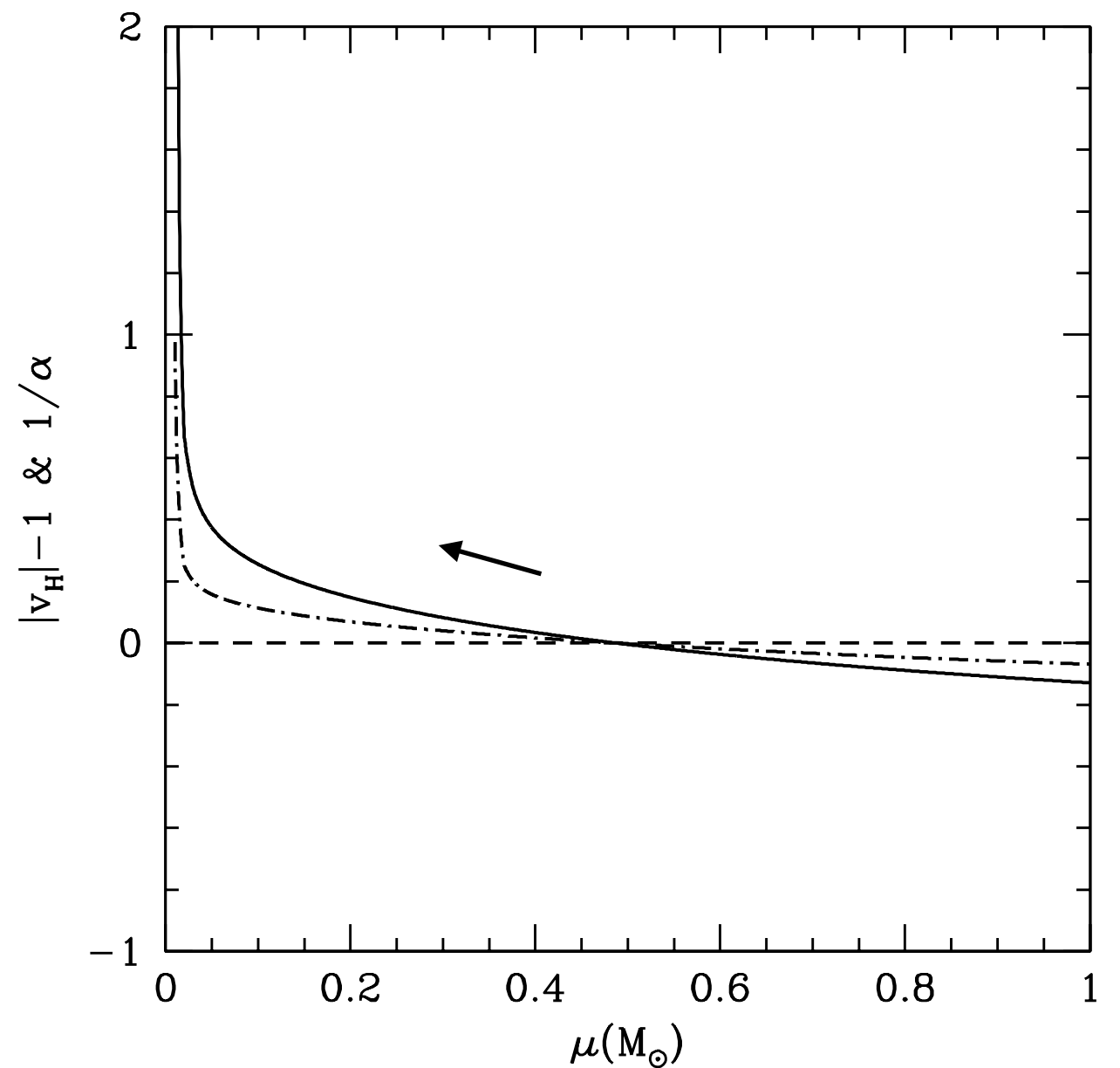
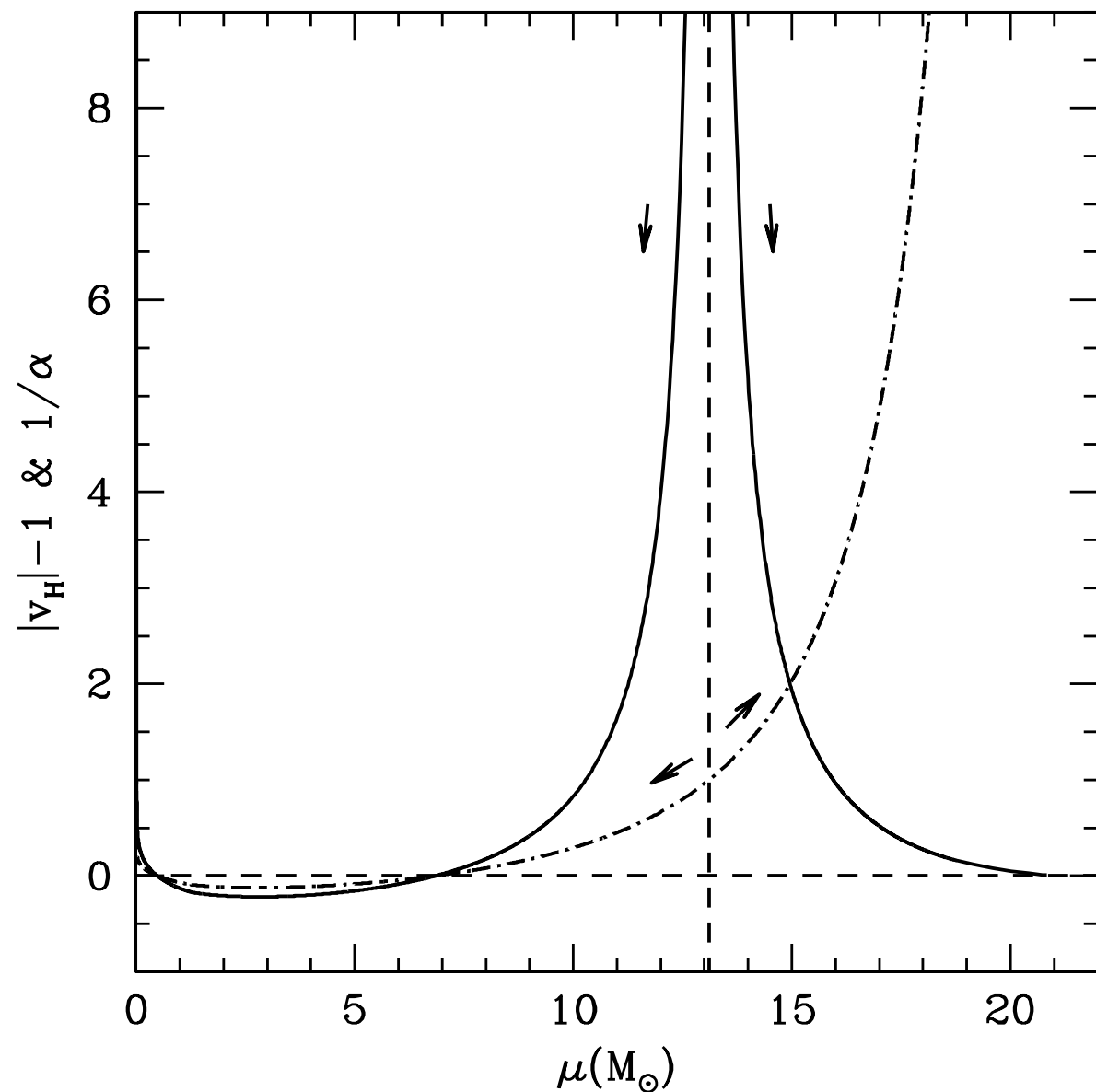
General scheme for in/out-going horizon evolution



$$p = K\rho^\gamma \quad (\gamma = 5/3, \text{HOM I.C.})$$

$$\alpha = \frac{4\pi R_H^2 (e + p)}{1 - 4\pi R_H^2 (e - p)}$$

$$v_H = \frac{1 + 8\pi R_H^2 p}{1 - 8\pi R_H^2 e}$$



Oppenheimer-Snyder collapse

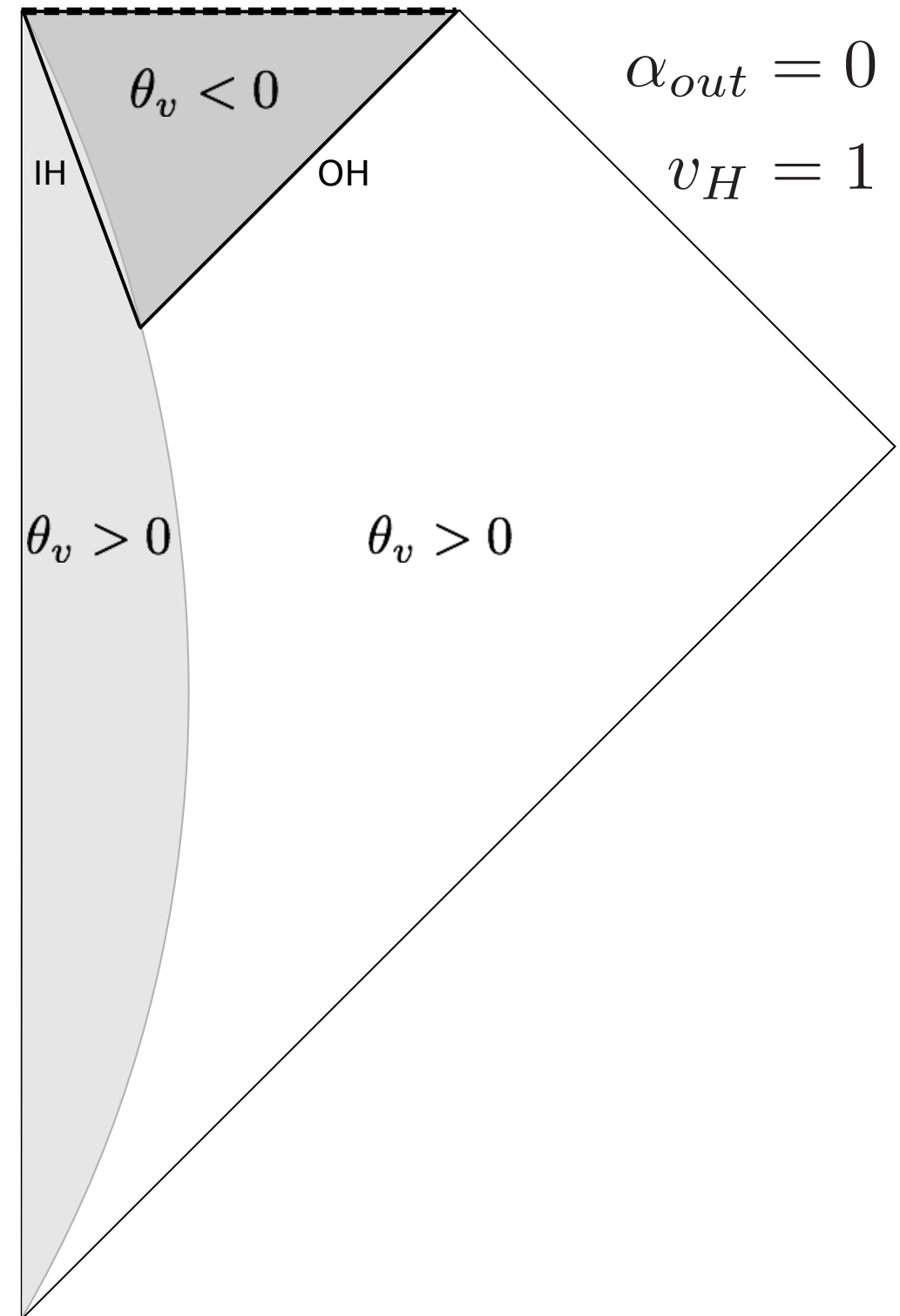
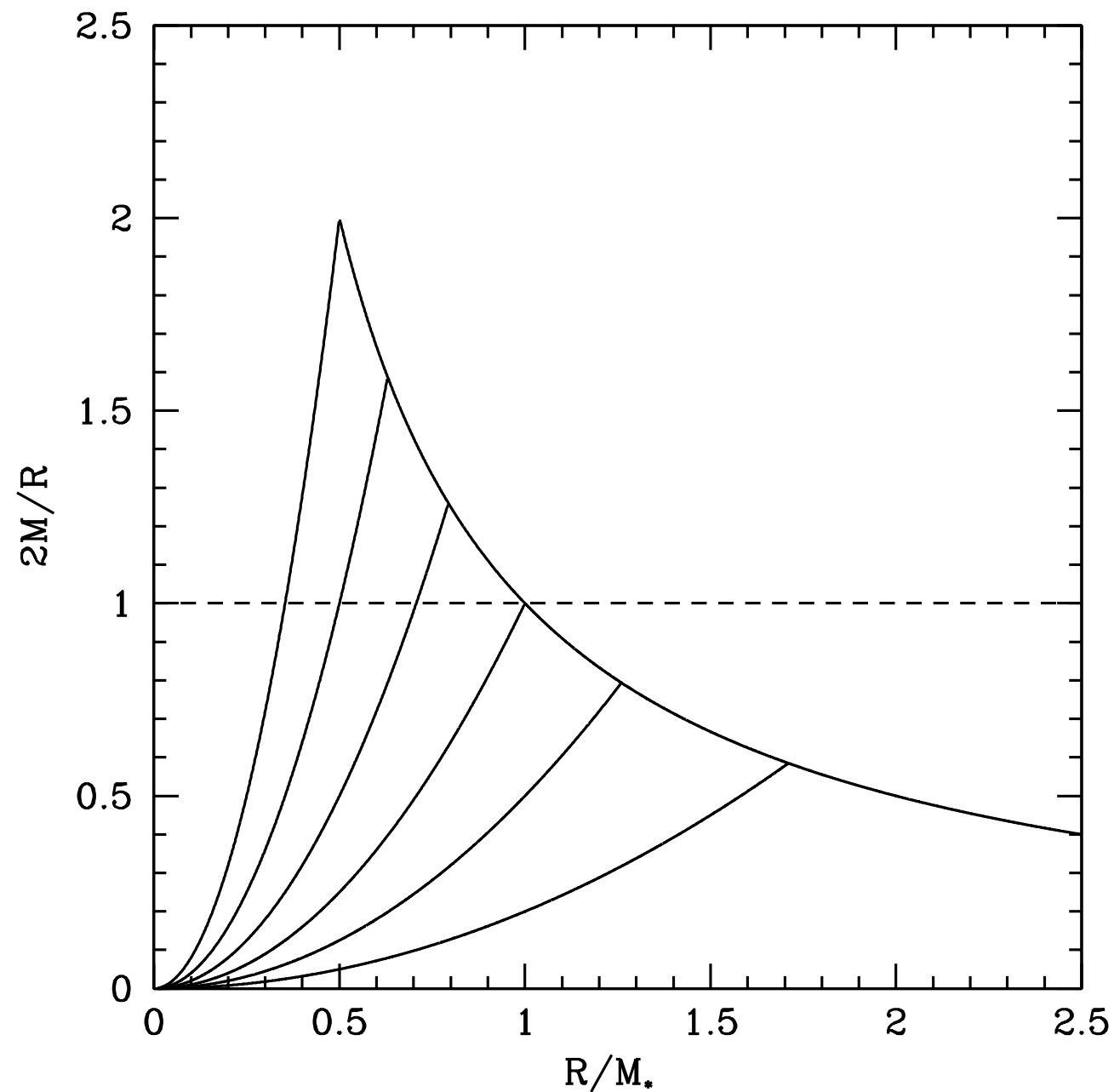
$$\frac{2M}{R} = \frac{8}{3}\pi R^2 e$$

$$\alpha_{in} = -3$$

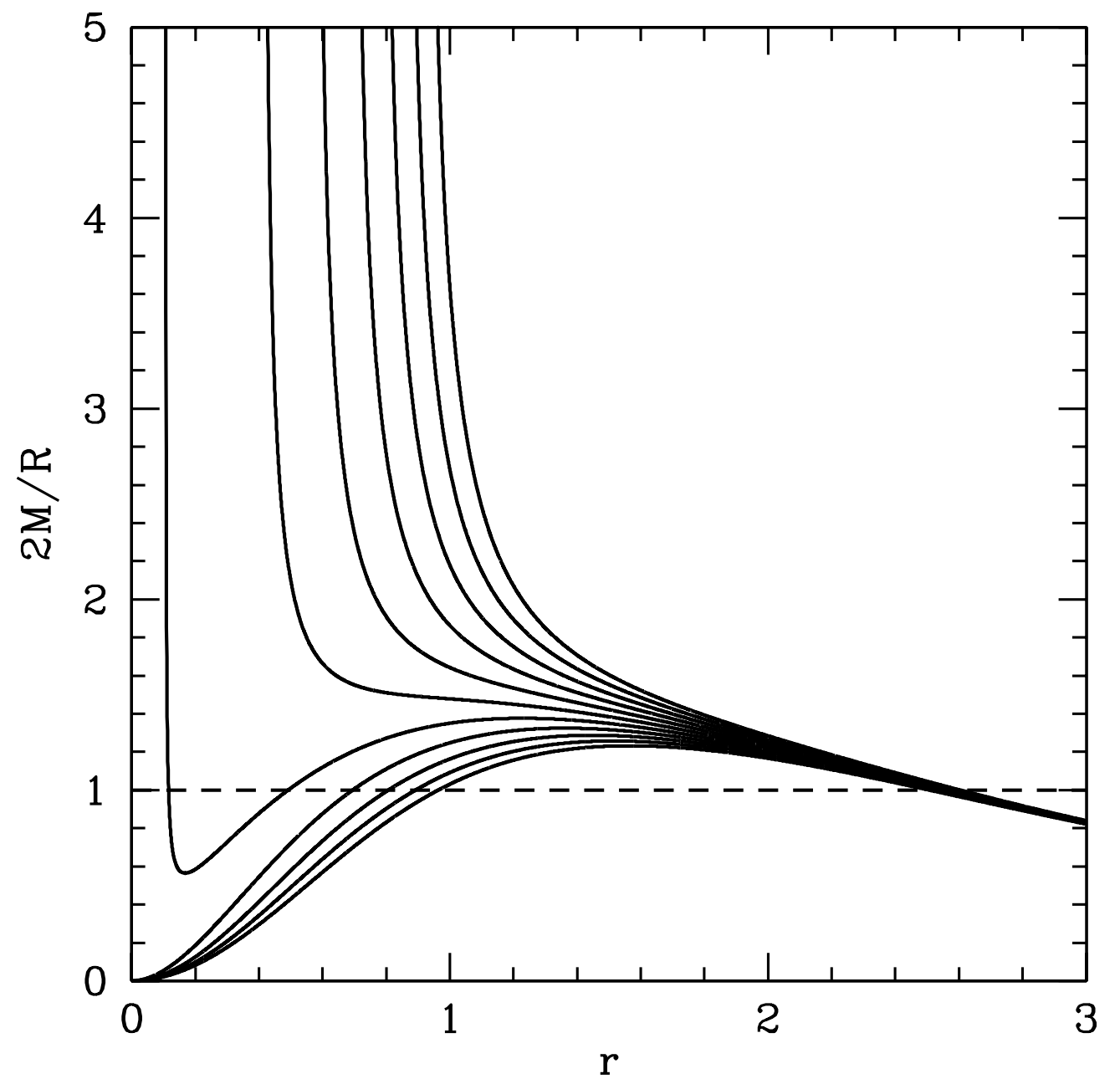
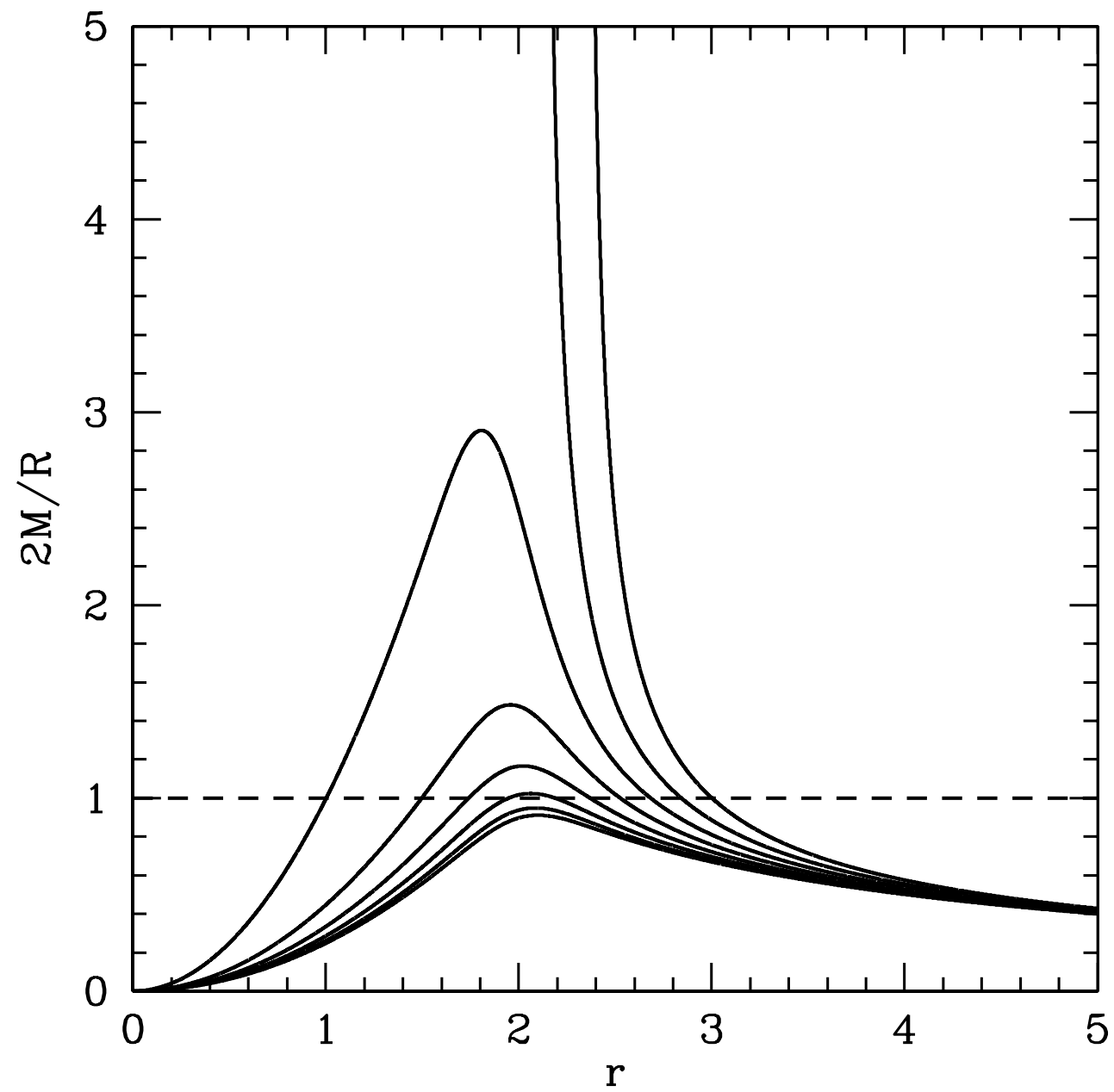
$$v_H = -1/2$$

$$\alpha_{out} = 0$$

$$v_H = 1$$

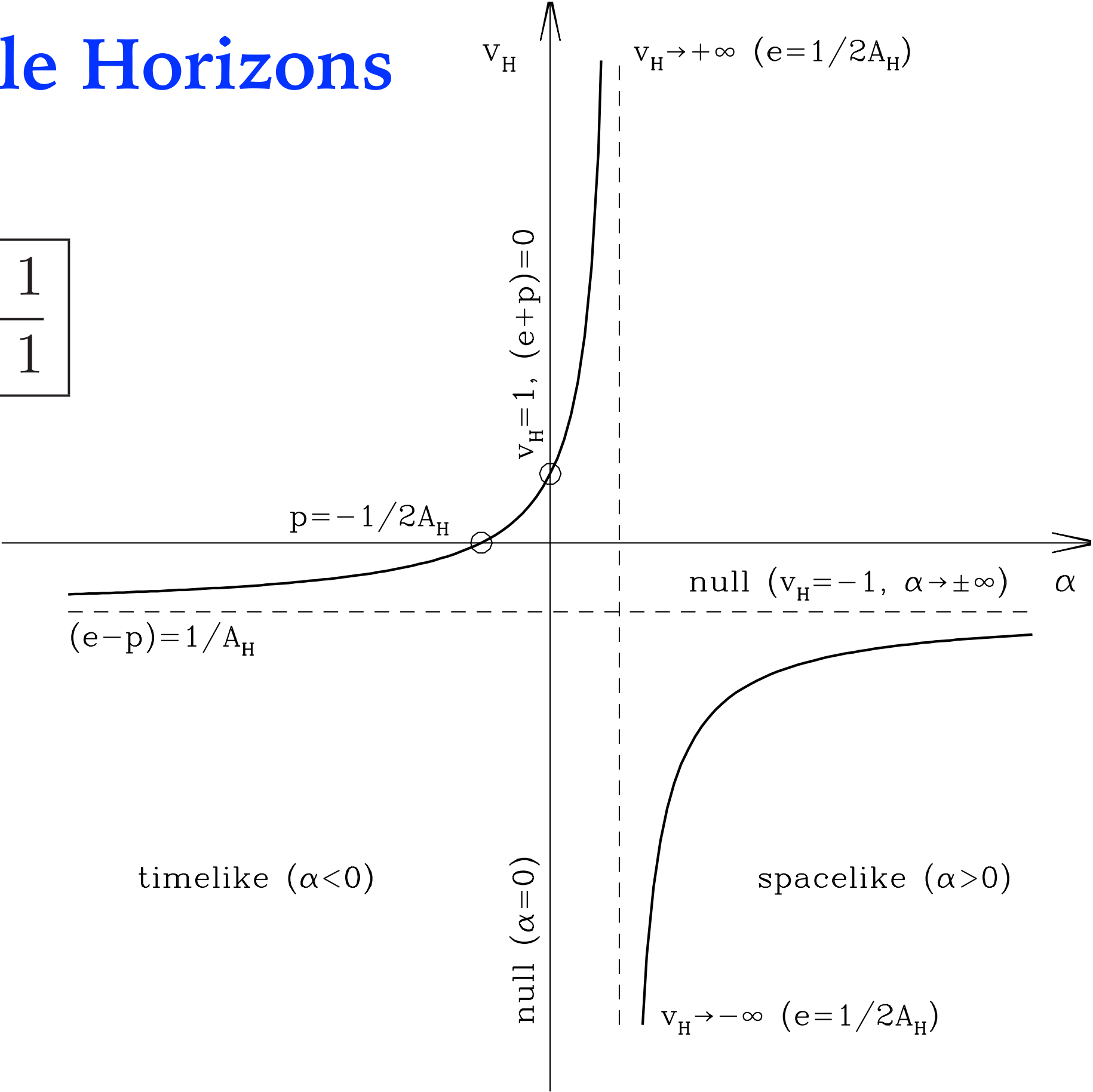


LTB collapse (zero pressure)



Black Hole Horizons

$$v_H = -\frac{\alpha + 1}{\alpha - 1}$$

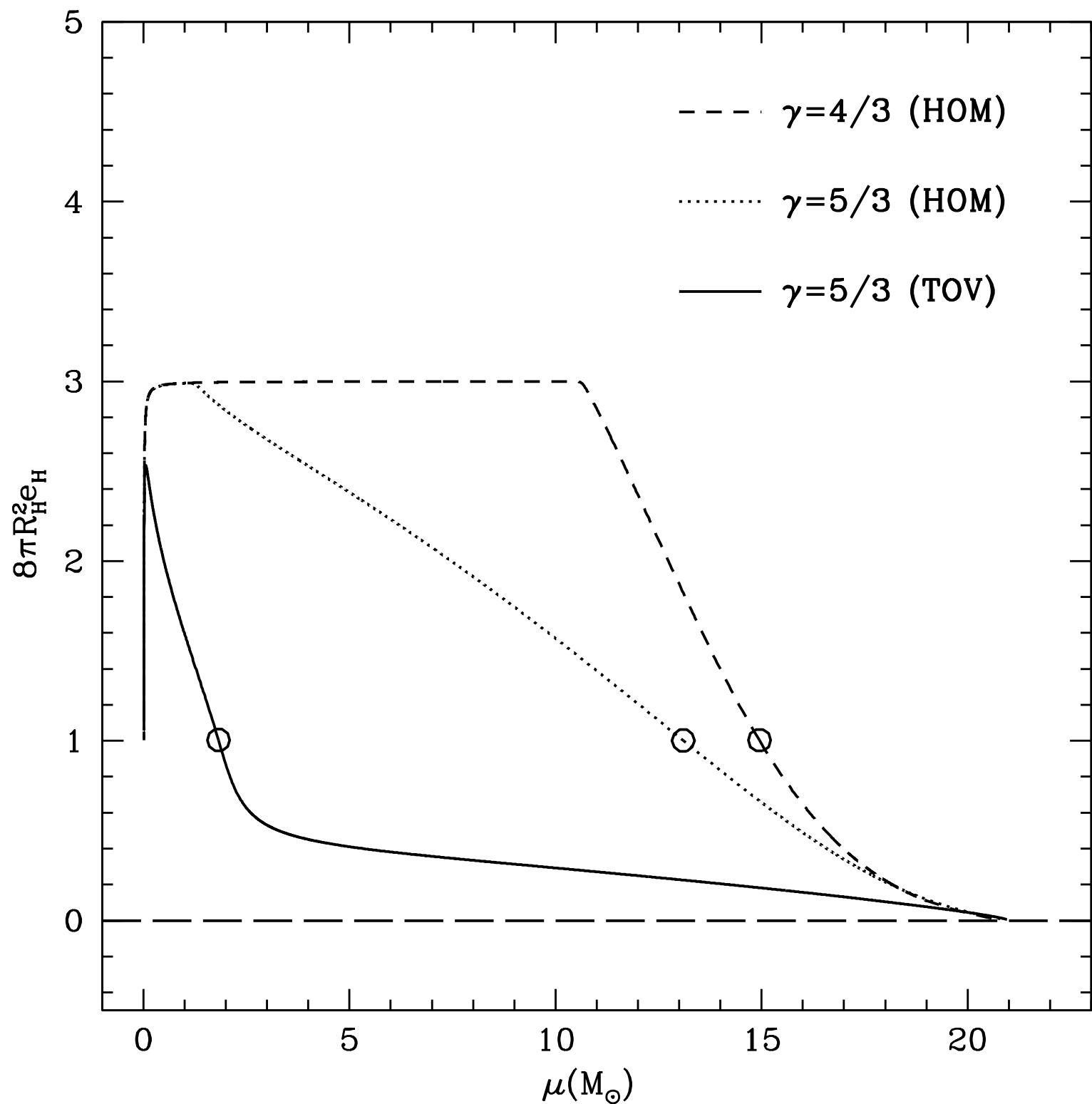


Causal Nature Summary

$$\alpha = \frac{4\pi R_H^2 (e + p)}{1 - 4\pi R_H^2 (e - p)}$$

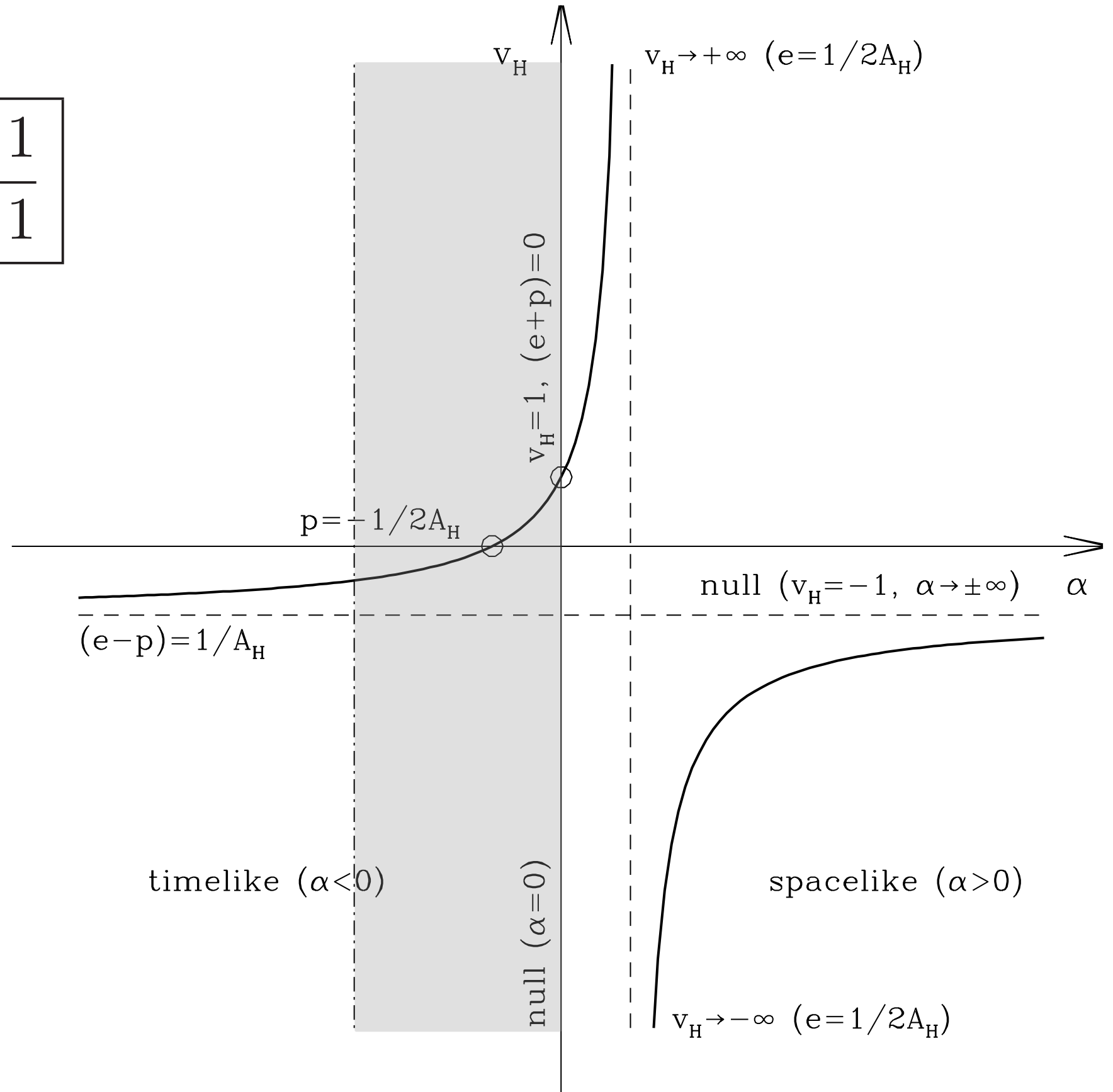
$$v_H = \frac{1 + 8\pi R_H^2 p}{1 - 8\pi R_H^2 e}$$

$$(\alpha = 1) \Rightarrow e = \frac{1}{2A_H}$$



Black Hole Horizons (quantum region)

$$v_H = -\frac{\alpha + 1}{\alpha - 1}$$



Conclusions & Future perspectives

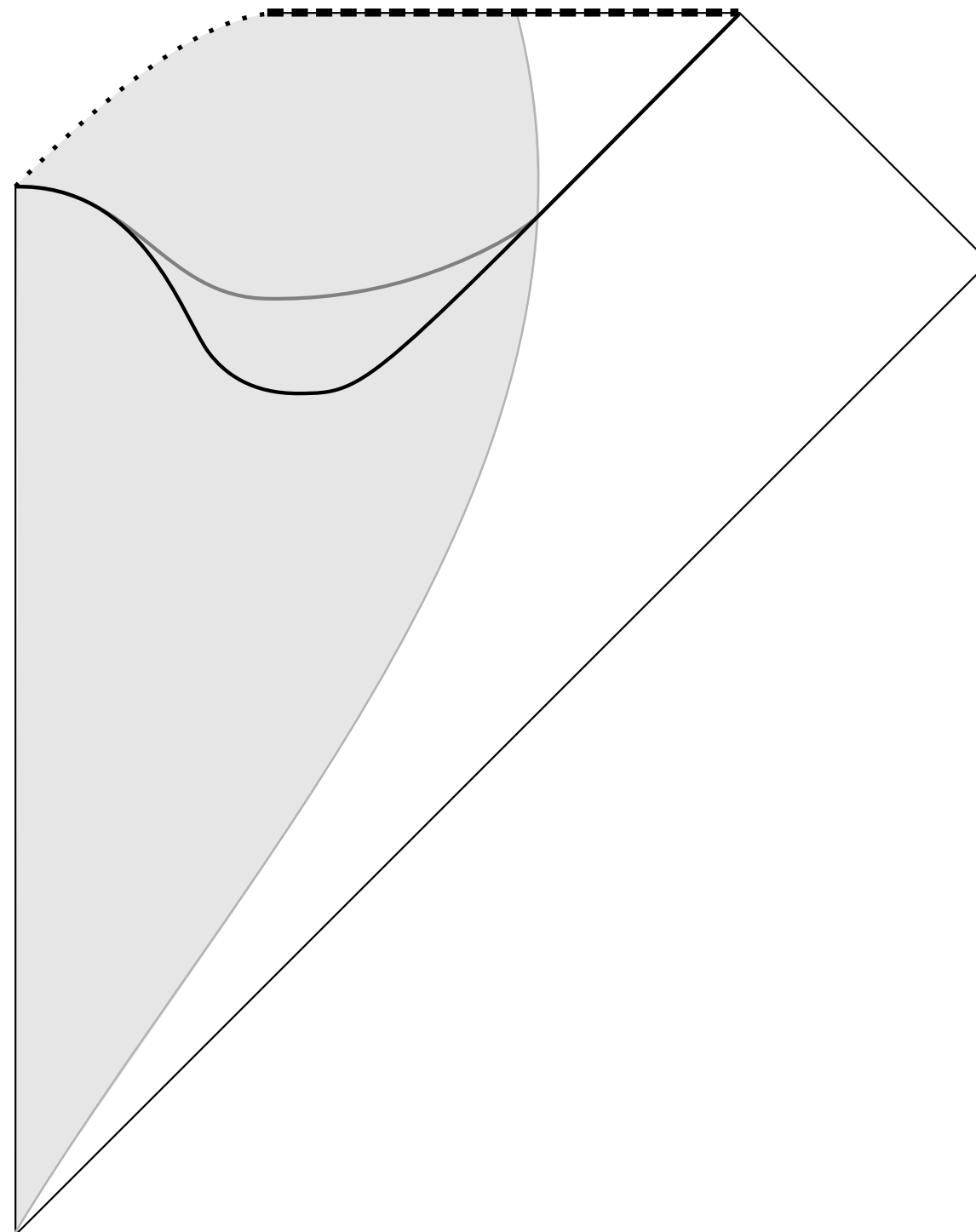
- With the Misner-Sharp equations (cosmic time slicing) we have studied the causal nature of trapping horizons appearing in gravitational collapse for polytropic stars forming black holes using a spherically symmetric Lagrangian numerical code.
- Within the classical regime of GR we have observed space-like outgoing horizons and space-like/time-like ingoing horizons depending on the choice of the equations of state and initial conditions. **Pressure seems to play a key role!**
- The conditions of **horizon formation** and **annihilation** are independent of the initial conditions:

$$\alpha = 1, \quad v_H = \pm\infty$$

- The formalism developed seems to show the possibility of incorporating quantum effects within the classical formulation of the GR-hydro equations modifying the equation of state accordingly to quantum gravity.

Can we get a bounce instead of a singularity?

Penrose Diagram of Gravitational Collapse with pressure



$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$

$$T_{\mu\nu} = (e + p)u_\mu u_\nu - pg_{\mu\nu}$$

COSMIC TIME

$$D_t \equiv \frac{1}{a} \left(\frac{\partial}{\partial t} \right) \quad D_r \equiv \frac{1}{b} \left(\frac{\partial}{\partial r} \right)$$

$$U \equiv D_t R \quad \Gamma \equiv D_r R$$

$$D_t U = - \left[\frac{\Gamma}{(e + p)} D_r p + \frac{M}{R^2} + 4\pi R p \right]$$

$$D_t \rho = - \frac{\rho}{\Gamma R^2} D_r (R^2 U)$$

$$D_t e = \frac{e + p}{\rho} D_t \rho$$

$$D_t M = -4\pi R^2 p U$$

$$D_r a = - \frac{a}{e + p} D_r p$$

$$D_r M = 4\pi R^2 \Gamma e$$

$$\Gamma^2 = 1 + U^2 - \frac{2M}{R}$$

$$ds^2 = -a^2 dt^2 + b^2 dr^2 + R^2 d\Omega^2$$

- Proper time / space derivative

- 4-velocity & Lorentz factor

- Euler equation

- Continuity equation

- Mass conservation

$$dM = aU dt + b\Gamma dr$$

- Lapse equation / pressure gradients

- Constraint equation

Equation of State

energy density: $e = \rho(1 + \epsilon)$

pressure: $p = (\gamma - 1)\rho\epsilon$

rest mass density

adiabatic index - particle degree of freedom

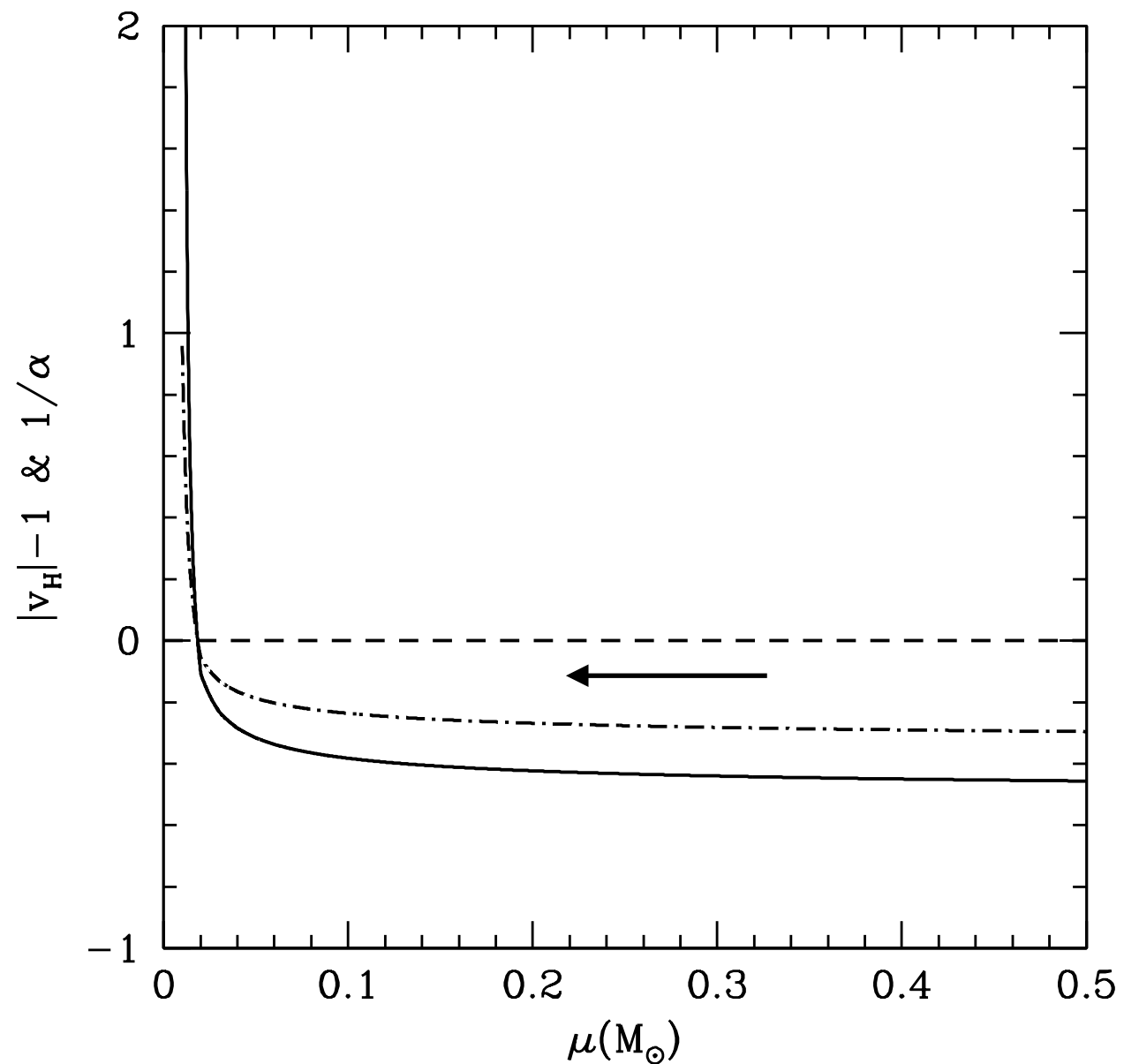
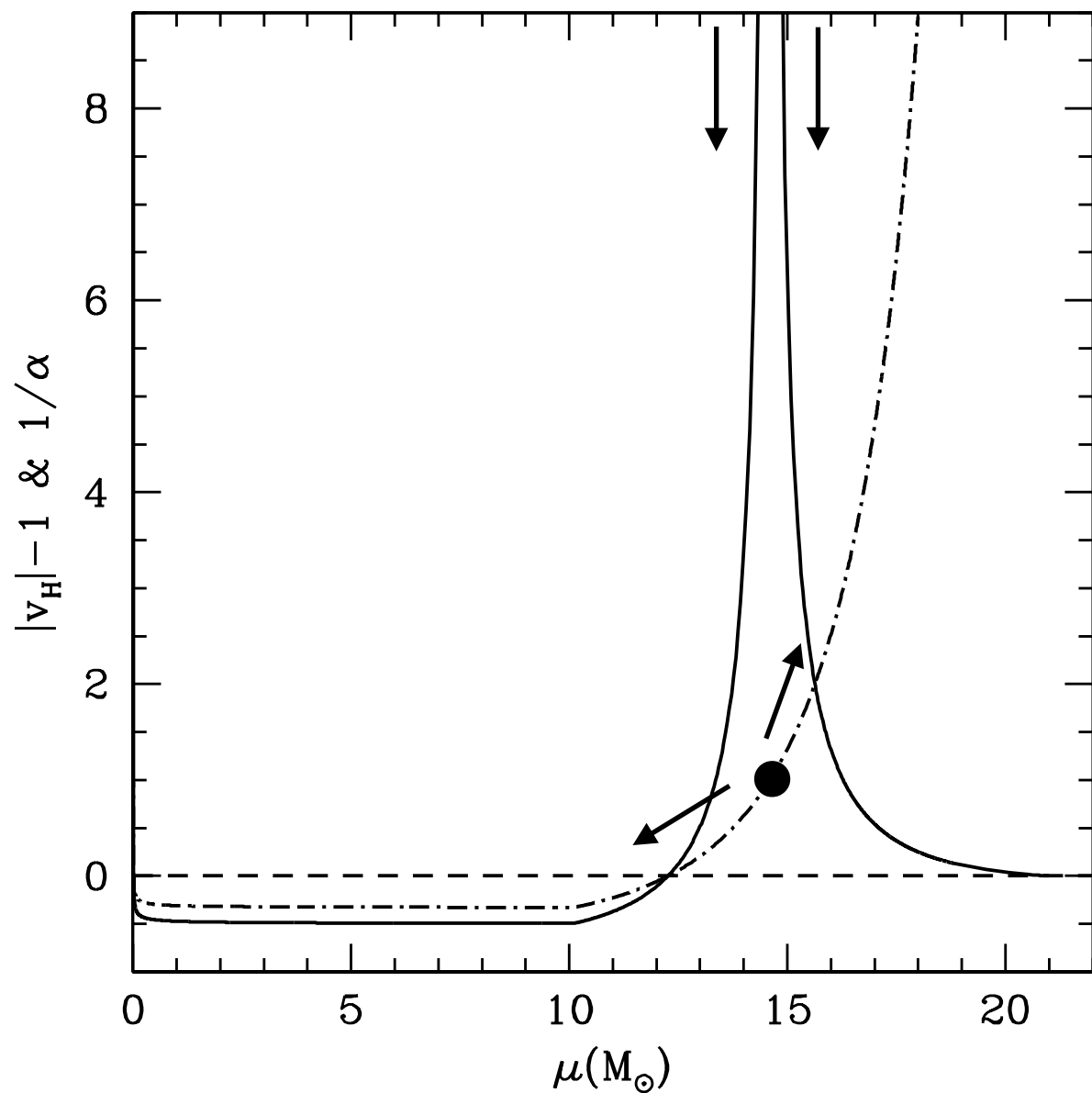
specific internal energy (velocity dispersion)

- Barotropic fluid (no rest mass density): $p = we$ with $w \in [0, 1]$
 - radiation dominated era: $w = 1/3$ RADIATION ($\gamma = 4/3$)
 - matter dominated era: $w = 0$ DUST ($\gamma = 1$)
- Polytropic fluid: $p = K(s)\rho^\gamma$ ($\gamma = 5/3, 4/3, 2$)
 - If the fluid is adiabatic (no entropy change): $K(s) = K$ (constant)

$$p = K\rho^\gamma \quad (\gamma = 4/3, \text{HOM I.C.})$$

$$\alpha = \frac{4\pi R_H^2(e + p)}{1 - 4\pi R_H^2(e - p)}$$

$$v_H = \frac{1 + 8\pi R_H^2 p}{1 - 8\pi R_H^2 e}$$



$$p = K\rho^\gamma \quad (\gamma = 5/3, \text{TOV I.C.})$$

$$\alpha = \frac{4\pi R_H^2 (e + p)}{1 - 4\pi R_H^2 (e - p)}$$

$$v_H = \frac{1 + 8\pi R_H^2 p}{1 - 8\pi R_H^2 e}$$

