

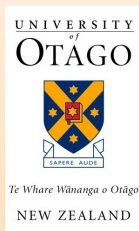
# The conformally invariant wave equation near the cylinder at spacelike infinity

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# Introduction

- Conformal compactification of asymptotically simple spacetimes: Mathematical difficulties can arise near **spacelike infinity**  $i^0$ .
- Helmut Friedrich's **generalised conformal field equations** are adapted to an appropriate treatment of  $i^0$ , which was thought of before as a point, but is now **'blown up' to a cylinder**  $S^2 \times [-1, 1]$  that connects past and future null infinity  $\mathcal{I}^\pm$ .
- To what extent can fields near  $i^0$  be obtained numerically?
- *Toy model: Conformally invariant wave equation* (in 4 dimensions) on Minkowski or Schwarzschild background,

$$g^{ab} \nabla_a \nabla_b f - \frac{R}{6} f = 0.$$

Conformal weight is  $-1$ , i.e. if  $\tilde{f}$  solves the equation for a metric  $\tilde{g}_{ab}$ , then  $f = \Theta^{-1} \tilde{f}$  solves the equation for  $g_{ab} = \Theta^2 \tilde{g}_{ab}$ .

- Typical problem in some numerical methods: Time step restricted by **CFL condition**.  $\rightarrow \mathcal{I}^+$  cannot be reached with a finite number of steps. Here: **fully pseudospectral time evolution**

# The fully pseudospectral scheme

- *Basic idea for solving 1 + 1 dimensional problems:*

[JH and Ansorg 2009]

- 1 Map the physical domain to one (or several) unit square(s) by introducing **spectral coordinates**  $(\sigma, \tau) \in [0, 1] \times [0, 1]$  such that the surface on which initial data are given corresponds to  $\tau = 0$ .
- 2 Enforce the **initial conditions** by expressing the unknown function  $f(\sigma, \tau)$  in terms of another unknown  $f_2(\sigma, \tau)$  via

$$f(\sigma, \tau) = \begin{cases} f_0(\sigma) + \tau f_1(\sigma) + \tau^2 f_2(\sigma, \tau), & \text{if function values and first} \\ & \text{derivatives are given} \\ f_0(\sigma) + \tau f_2(\sigma, \tau), & \text{if only function values are} \\ & \text{given} \end{cases}$$

depending on the type of problem (first/second order equations, characteristic/Cauchy initial value problem).

# The fully pseudospectral scheme

- 3 Approximate the new unknown  $f_2$  via

$$f_2(\sigma, \tau) \approx \sum_{j=0}^{N_\sigma} \sum_{k=0}^{N_\tau} c_{jk} T_j(2\sigma - 1) T_k(2\tau - 1)$$

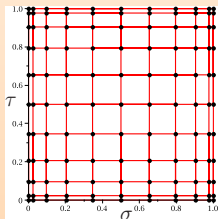
for given spectral resolution (number of polynomials)

$n_\sigma \equiv N_\sigma + 1$  and  $n_\tau \equiv N_\tau + 1$ .

- 4 Obtain an algebraic system of equations by requiring that the equation(s), and suitable boundary or regularity conditions, are satisfied at a set of collocation points. We choose **Gauss-Lobatto nodes**  $(\sigma_j, \tau_k), j = 0, \dots, N_\sigma, k = 0, \dots, N_\tau$ :

$$\sigma_j = \sin^2 \left( \frac{\pi j}{2N_\sigma} \right), \quad \tau_k = \sin^2 \left( \frac{\pi k}{2N_\tau} \right)$$

- 5 Starting from some initial guess, solve this system iteratively with the **Newton-Raphson method**.



# Minkowski background

- **Minkowski metric**  $\tilde{g}$  in spherical coordinates

$$\tilde{g} = d\tilde{t}^2 - d\tilde{r}^2 - \tilde{r}^2(d\theta^2 + \sin^2\theta d\phi^2)$$

- Compactification:

$$\tilde{u} = \tilde{t} - \tilde{r}, \quad \tilde{v} = \tilde{t} + \tilde{r}; \quad \tilde{u} = \tan u, \quad \tilde{v} = \tan v, \quad T = v + u, \quad R = v - u$$

$$T = 2t \cos(r), \quad R = 2r$$

$$g = dt^2 - 2t \tan r dt dr - [1 + (1 - t^2) \tan^2 r] dr^2 - \sin^2 r (d\theta^2 + \sin^2 \theta d\phi^2)$$

$$\Theta = \frac{\cos(t \cos r - r) \cos(t \cos r + r)}{\cos r}, \quad \tilde{g} = \Theta^{-2} g$$

- *Conformally invariant wave equation:*

$$(1 - t^2 \sin^2 r) f_{,tt} - 2t \sin r \cos r f_{,tr} - \cos^2 r f_{,rr} \\ - t(2 + \cos^2 r) f_{,t} - 2 \frac{\cos^3 r}{\sin r} f_{,r} + \cos^2 r f = 0$$

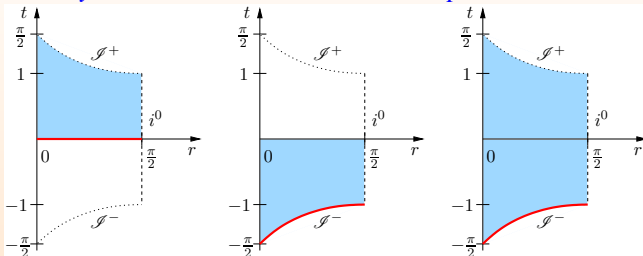
- *Regularity conditions:*

$$\text{at } r = 0: f_{,r} = 0, \quad \text{at } r = \frac{\pi}{2} (i^0): f_{,t} = 0$$

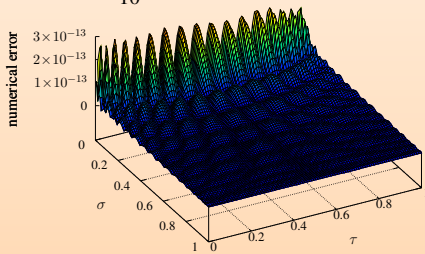
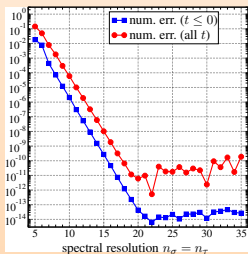
- *General (regular) solution:*  $f(t, r) = \frac{F(t \cos r - r) - F(t \cos r + r)}{\sin(r)}$

# Minkowski background

- *Cauchy or characteristic initial value problems:*



- *Example:* characteristic IVP,  $F = \frac{1}{10}x^3$  [Frauendiener and JH 2014]



# Schwarzschild background

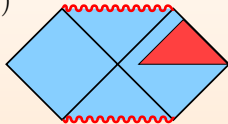
Solution in a neighbourhood of the cylinder at  $i^0$

- **Schwarzschild metric**  $\tilde{g}$  in isotropic coordinates

$$\tilde{g} = - \left( \frac{1 - \frac{2m}{\tilde{r}}}{1 + \frac{2m}{\tilde{r}}} \right) dt^2 - \left( 1 + \frac{m}{2\tilde{r}} \right)^4 (d\tilde{r}^2 + \tilde{r}^2 d\sigma^2)$$

- New coordinates  $0 < \rho < \rho_{\max} < 1$ ,  $0 \leq \tau < 1$   
[Friedrich, 2004]:

$$r = \frac{m}{2\tilde{r}}, \quad t = \frac{2\tilde{t}}{m}; \quad t = \int_r^\rho \frac{ds}{F(s)}, \quad r = \rho(1-\tau), \quad F(r) = \frac{r^2(1-r)}{(1+r)^3}$$



$$g = \frac{2\rho A}{r^2} d\rho d\tau - \frac{A}{r^2} [2(1-\tau) - A] d\rho^2 - d\sigma^2, \quad A := \frac{F(r)}{F(\rho)}, \quad \Theta = \frac{m(1+r)^2}{2r}$$

$\tau = 1$ :  $\mathcal{I}^+$ ,  $\rho = 0$ :  $i^0$ ,  $\rho = 1$ : coordinate singularity ( $A$  diverges)

- *Conformally invariant wave equation:*

$$[2(1-\tau) - A] f_{,\tau\tau} + 2\rho f_{,\tau\rho} - 2 \left[ 1 - \frac{\rho(1-2r)}{r(1-r^2)} A \right] f_{,\tau} + 4 \frac{\rho^2 A}{r(1+r)^2} f = 0$$

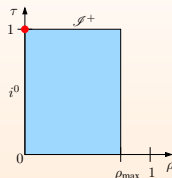
# Schwarzschild background

Solution in a neighbourhood of the cylinder at  $i^0$

- *General behaviour of  $f$  near  $i^0$ :*

Expansion:

$$f(\tau, \rho) = \phi_0(\tau) + \rho\phi_1(\tau) + \rho^2\phi_2(\tau) + \rho^3\phi_3(\tau) + \dots$$



Conformal wave eq. leads to ODEs for  $\phi_0, \phi_1, \dots$

- ①  $\phi_0(\tau) = c_0 + c_1 \ln \frac{1-\tau}{1+\tau} \rightarrow$  choose data with  $f_{,\tau}(0, 0) = -2c_1 = 0$
- ②  $\phi_1(\tau) = -2c_0\tau - \frac{2c_0+c_2}{1+\tau} + c_3 \rightarrow$  regular
- ③  $\phi_2(\tau) = \frac{4(2c_0-c_3)}{1+\tau)^2} (1-\tau)^2 \ln(1-\tau) +$  regular terms  
 $\rightarrow f_{,\rho\rho\tau\tau}$  singular, unless  $(f_{,\rho} + f_{,\tau\rho})(0, 0) = c_3 - 2c_0 = 0$
- ④  $\phi_3(\tau) = \frac{8(6c_0+5c_2+3c_4)}{3(1+\tau)^2} (1-\tau)^3 \ln(1-\tau) +$  regular terms  
 $\rightarrow f_{,\rho\rho\rho\tau\tau\tau}$  singular, unless  $(3f_{,\rho\rho} + 14f_{,\rho} - 36f)(0, 0) = 0$
- ⋮

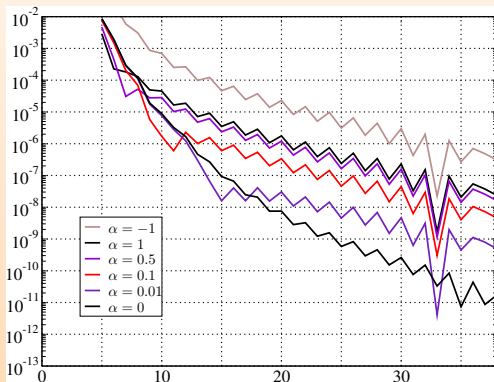


# Schwarzschild background

Solution in a neighbourhood of the cylinder at  $i^0$

*Family of initial data at  $\tau = 0$ :  $f = 2 + \sin(5\rho)$ ,  $f_{,\tau} = (\alpha - 5)\rho$*

- Choose  $\rho_{\max} = 0.5$
- For all  $\alpha$ :  $\phi_0$  regular
- For  $\alpha = 0$ :  $\phi_2$  regular



→ algebraic convergence

# Schwarzschild background

Including the horizon

- Again starting from isotropic coordinates, perform the following, modified coordinate transformation (with parameter  $d > \frac{1}{2}$ ):

$$t = \int_r^{\rho(2-\rho)} \frac{ds}{F(s)}, \quad r = \frac{\rho(1-\tau)}{w+\rho(1-\tau)}; \quad w = \frac{1-\rho}{(2-\rho)(1+d\rho)}, \quad F(r) = \frac{r^2(1-r)}{(1+r)^3}$$

- *Boundaries:*  $\mathcal{I}^+$  :  $\tau = 1$ ,  $i^0$  :  $\rho = 0$ , **horizon  $\mathcal{H}$  :  $\rho = 1$**
- Similar behaviour at the cylinder (conditions on initial data can remove the leading-order logarithmic singularities at  $\rho = 0$ ,  $\tau = 1$ )

# Schwarzschild background

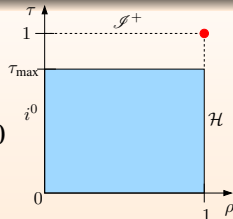
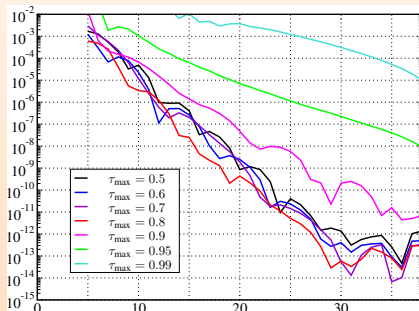
Including the horizon

- Numerical example:*

Domain:  $0 \leq \rho \leq 1, 0 \leq \tau_{\max} < 1$

Parameter:  $d = 2$

Initial data:  $f = 2 + \sin(5\rho), f_{,\tau} = -5\rho$  at  $\tau = 0$



→ spectral convergence and highly accurate solutions, if the numerical domain is not too close to the singularity at  $\rho = 1, \tau = 1$   
 [Frauendiener and JH, in preparation]

# Summary

- *Minkowski background:*

- For initial data subject to regularity conditions (at the intersection of the initial surface with  $i^0$  and at the origin) the conformally invariant wave equation has solutions that are regular up to  $\mathcal{I}^+$ .
- We can solve **Cauchy problems** in a neighbourhood of the cylinder or globally, and **characteristic initial value problems** with data at  $\mathcal{I}^-$ .  
→ **spectral convergence and highly accurate solutions**

- *Schwarzschild background:*

- Solutions generally have **logarithmic singularities** at the intersection of  $\mathcal{I}^+$  and  $i^0$ . The leading singularities can be removed with appropriate restrictions for the initial data.  
→ **algebraic convergence, still very accurate solutions**
- We can use coordinates that include the event horizon  $\mathcal{H}$  and solve the wave equation in a domain with  $0 \leq \tau \leq \tau_{\max} < 1$  avoiding  $\mathcal{I}^+$ .  
→ **spectral convergence, if  $\tau_{\max}$  is not too close to 1**